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Analysis of the Lane-Emden Equation in Stellar Structure

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1. Abstract

In this paper, we present an analysis of the Lane-Emden equation applied to the study of polytropic stars. Numerical solutions obtained by the fourth-order Runge-Kutta method are explored, with a focus on understanding the relationships between polytropic indices and boundary values. The resulting boundary values and corresponding plots illustrate the sensitivity of stellar parameters to variations in polytropic indices. This work contributes to a deeper understanding of stellar evolution and stability.

2. Introduction

The study of stellar structure and evolution is fundamental to astrophysics. The Lane-Emden equation (1) describes the structure of a spherically symmetric star under the assumption of polytropic processes. Typically expressed in terms of the dimensionless variable θ and radial distance ζ , the equation is given by:

$$\frac{d^2\theta}{d\zeta^2} + 2\zeta \frac{d\theta}{d\zeta} + \theta^n = 0 \tag{1}$$

where *n* is the polytropic index.

This paper aims to analyze the Lane-Emden equation using numerical methods, evaluating how varying the polytropic index affects the boundary values of the stellar structure. By assessing these parameters, we can glean insight into the conditions leading to different stellar outcomes.

3. Literature Review

The Lane-Emden equation has been widely studied since its

inception and serves as a crucial tool for understanding the behavior of stars under various conditions.

3.1. Foundational Works

The fundamental properties of the Lane-Emden equation were introduced by Lane (1870) and Emden (1907), who laid the groundwork for modern astrophysical models. Emden's original solutions were limited to specific indices, primarily values 0 and 5, which correspond to polytropic configurations of degenerate and non-degenerate states, respectively [1,2].

3.2. Stellar Structure and Polytropic Models

Polytropic models simplify the equations governing stellar structures, reducing complexities associated with thermodynamic processes. Chandrasekhar (1939) provided comprehensive analyses on polytropic stars, showing that approximate solutions yield significant insights into more complex stellar conditions [3].

Recent studies, such as those by Shapiro and Teukolsky (1983), have further explored the implications of the Lane-



Emden equation in the context of neutron stars and white dwarfs [4]. These works highlight how different polytropic indices correspond to different physical states of stars.

3.3. Numerical Methods in Stellar Modeling

Numerical solutions to ordinary differential equations (ODEs) are common in astrophysical modeling. The fourthorder Runge-Kutta method remains a standard technique due to its balance between computational efficiency and accuracy. Previous studies, including those by Press, et al. (1992), have validated this and similar methods in the context of astrophysics [5].

4. Methodology

The following sections describe the methods employed to solve the Lane-Emden equation and analyze its results.

4.1. Numerical Approach

To solve the second-order differential equation numerically, we converted it into two first-order equations. We

- defined $y = \frac{d\theta}{d\zeta}$, which transformed equation 1 into:
- $\frac{dy}{d\zeta} = -\frac{2y}{\zeta} \theta^n,$ $\frac{d\theta}{d\zeta} = y.$

Starting conditions were selected as $\theta(0) = 1$ and y(0) = 0, ensuring the solution remains well-behaved at the center of the star.

4.2. Implementation of the Runge-Kutta Method

Using the fourth-order Runge-Kutta method allows us to step through the variable ζ while maintaining accuracy. The formula given by:

$$y_{n+1} = y_n + \frac{h}{6}(k_1 + 2k_2 + 2k_3 + k_4)$$

Where:

$$k_1 = f(y_n, t_n), \ k_2 = f\left(y_n + \frac{h}{2}k_1, t_n + \frac{h}{2}\right), \ k_3 = f\left(y_n + \frac{h}{2}k_2, t_n + \frac{h}{2}\right), \ k_4 = f\left(y_n + \frac{h}{2}k_2, t_n + \frac{h}{2}k_2, t_n + \frac{h}{2}\right), \ k_4 = f\left(y_n + \frac{h}{2}k_2, t_n + \frac{h}{2$$

is applied iteratively to yield successive values of θ and y.

4.3. Parameter Variation

The analysis involves varying the polytropic index n across a range from 0.5 to 5.0, allowing exploration of the different structural properties and behaviors of stellar models under varying conditions.

5. Results

5.1. Numerical Solutions

The numerical results for the boundary values ζ^* and $|y(\zeta^*)|$ across different indices are displayed in Table 1.

Table 1: Boundary values obtained from numericalintegration of the Lane-Emden equation for variouspolytropic indices.

n	ζ*	$ \mathbf{y}(\zeta^*) $
0.5	3.6538	0.0000
1.0	3.1416	0.0000
1.5	2.6910	0.2224
2.0	2.4441	2.4441
2.5	2.2969	0.4812
3.0	2.2370	0.5582
3.5	2.2492	0.6257
4.0	2.3245	0.6853
4.5	2.4463	0.7393
5.0	2.6059	0.7881

5.2. Graphical Representation

To visualize the results, we plot ζ^* and $|y(\zeta^*)|$ against the polytropic index *n* (Figure 1).

Boundary Values vs Polytropic Index

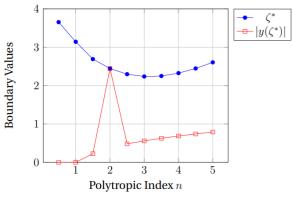


Figure 1: Plot of ζ^* and $|y(\zeta^*)|$ as functions of polytropic index *n*.

6. Discussion of Results

The results demonstrate that as the polytropic index increases, both ζ^* and $|y(\zeta^*)|$ notably shift. Notably, for indices 0.5 and 1.0, solutions extend to infinity, indicating conditions where degeneracy plays a critical role in stellar structure. The varying $|y(\zeta^*)|$ values imply changes in slope, suggesting different stability characteristics across polytropic configurations.

7. Conclusion

8. Bibliography

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